## ON THE NATURE OF PERIODICALLY PULSATING RADIATION SOURCES

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A change in the character of maser generation in a two-level system is found when the initial population inversion exceeds some threshold value proportional to the square root of the total number of atoms. Above this threshold, the number of photons begins to grow exponentially with time and the pulse with short leading edge and broadened trailing edge is generated. In this work, we attempt to explain the nature of this threshold. Coherent pulse duration estimated by its half-width increases significantly with increasing inversion, if all other parameters are fixed and the absorption is neglected. The inclusion of the energy loss of photons leads to the fact that the duration of coherent pulse is almost constant with increasing inversion, at least well away from the threshold. If there is exist a recovery mechanism for the population inversion, the pulsating mode of stimulated emission generation become possible. The integral radiation intensity at this may be increased several times. This approach can be used for analysis of the cosmic radiation that might help explain a great variety of pulsating radiation sources in space.

## I. INTRODUCTION

Description of physical phenomena based on the systems of partial differential equations, derived from the observations and experimental facts, often conceals from an investigator some essential features, especially in those cases, when the researchers do not expect to find anomalies and 
qualitative changes in the dynamics of systems in given 
range of variables and parameters. Namely such a case of 
unusual behavior of a two-level quantum system was found 
in attempting to separate a stimulated component from the 
total radiation flow.

In the beginning of the past century, A. Einstein has proposed the model of two-level system, which has demonstrated the possibility of generation of both spontaneous and induced (stimulated) emission when the initial population inversion is sufficiently large [1]. Usually, the term spontaneous emission denotes the emission of oscillator (or other emitter) which not forced by external field of the same frequency. As for other influences on the characteristics of the spontaneous emission, there is nothing to say definitely. Although the dynamics of spontaneous processes usually shows a steady recurrence and invariance, there is evident [2] that the characteristics of the spontaneous processes can vary with change of environment. By induced or simulated emission is usually meant the emission produced because of an external field action on the emitting source at the radiation frequency.

There were difficulties in the quantum description with interpretation of the stimulated emission as coherent, where in contrast to the classical case it was impossible to say anything about the phases of the fields emitted by individual atoms and molecules. However, C. Townes believed [3] that "the energy delivered by the molecular systems has the same field distribution and frequency as the stimulating radiation and hence a constant (possibly zero) phase difference"

If we assume, relying upon the results of the studies of fluctuation correlations in the laser radiation [4, 5], that a stimulated emission has a high proportion of the coherent component, one can find a threshold of coherent radiation at a certain critical value of population inversion [6]. The specific feature of this threshold is that it follows from the condition that the initial value of the population inversion is equal to the square root of the total number of states. On the other hand, the change in the nature of the process near the threshold is evident, even without making any other assumptions. Above this threshold, the number of photons begins to grow exponentially with time. Herewith, below the threshold there no exponential growth.

It is known that at low levels of spontaneous component and far above the maser generation threshold the number of photons growths exponentially and the radiation is largely a coherent [7, 8]. The meaningful indicator of the collective character of stimulated emission is the so-called photon degeneracy, which is defined as the average photon number contained in a single mode of optical field (see, for example [9]). For the incoherent light, this parameter does not exceed unity, but for even the simplest He-Ne maser it reaches the value of  $10^{12}$  as was shown in the early works (see [71).

It is of interest to go further and analyze the consequences of consideration of the spontaneous emission as a random process (at least, in a homogeneous medium) and induced process as a coherent process. It is clear that the separation of total radiation into two category: the stimu-

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lated – coherent and spontaneous – random or incoherent will be idealized simplification. However, such separation may explain, at least qualitatively, the nature of the radiation emitted by two-level quantum system near to exposed threshold.

Another indirect proof of the existence of such a threshold is the following observation. The intensity of the spontaneous emission, which is non-synchronized (randomly distributed) over oscillators phases is known to be proportional to their number. The intensity of the coherent stimulated emission is proportional in turn to the square of the number of oscillators. It is easy to see that the exposed threshold corresponds to the case when the intensity of spontaneous and stimulated coherent radiation become equal.

In papers [6, 10], we have shown that under these conditions the pulse of stimulated radiation with a characteristic profile is formed when the initial population inversion slightly exceeds the threshold. The leading edge of the pulse due to the exponential growth of the field is very sharp due to the exponential growth of the field, and the trailing edge is rather broadened. Further overriding of the threshold, that is growing of the initial population inversion, results in the ratio of the trailing edge duration to the leading edge duration becomes greater. At large times the incoherent radiation dominates.

Because very small value of the initial population inversion can provide generation of pulses of stimulated radiation, it is of interest to determine the shape of these pulses for different values of the initial population inversion levels and when the field energy absorption should be taken into account. These pulses can be easily detected in experiments. In addition, after experimental validation of this model, it will be possible to use these approaches for analysis of the cosmic radiation that might help explain such abundance of coherent radiation sources in space.

In this paper, we study the characteristics of the pulses of stimulated (coherent) radiation as a function of the initial inversion and absorption level in the system. The dynamics of the emission process in the simplified model is compared with the dynamics of change in the number of quanta in the traditional model, where the separation into spontaneous and stimulated components is not carried out.

If there is exist a recovery mechanism for the population inversion, the pulsating mode of stimulated emission generation becomes possible. The integral radiation intensity at this may be increased several times. This approach can be used for analysis of the cosmic radiation that might help explain a great variety of pulsating radiation sources in space. In present work, we investigate the characteristics of the periodic pulse generation depending on the initial inversion, the pumping level and the absorption rate.

- [1] A. Einstein, Physikalische Gesellschaft Zürich 18 (1916).
- [2] P. Goy, J. Raimond, M. Gross, and S. Haroche, Phys. Rev. Lett. 50, 1903 (1983).
- [3] C. Townes, IEEE Spectrum, 2 (1965).
- [4] R. H. Brown and R. Q. Twiss, in Proceedings of the Royal Society of London A: Mathematical, Physical and Engineering Sciences, Vol. 242 (The Royal Society, 1957) pp. 300–324.
- [5] R. H. Brown and R. Twiss, in Proceedings of the Royal Society of London A: Mathematical, Physical and Engineering Sciences, Vol. 243 (The Royal Society, 1958) pp. 291–319.
- [6] V. Kuklin and A. Zagorodny, in XIV Khariton topical scientific readings" High-Power Pulsed Electrophysics", VNIIEF, Sarov,

Russia (2012).

- [7] G. Birnbaum and H. Malamud, Physics Today 18, 80 (1965).
- [8] N. Blotmbergen, Nonlinear Optics. A Lecture Note (WA Benjamin, Inc. New York-Amsterdam, 1965).
- [9] G. S. He and S. H. Liu, Physics of nonlinear optics, Vol. 216 (World Scientific, 1999).
- [10] A. Kirichok, V. Kuklin, A. Mischin, A. Pryjmak, and A. Zagorodny, Problems of Atomic Science and Technology, series: âĂIJPlasma Electronics and New Methods of AccelerationâĂİ, 267 (2013).
- [11] H. Statz and G. DeMars, in Quantum electronics (1960) p. 530.